Measurement of the Azimuthal Correlation between the most Forward Jet and the Scattered Positron in Deep Inelastic Scattering at HERA

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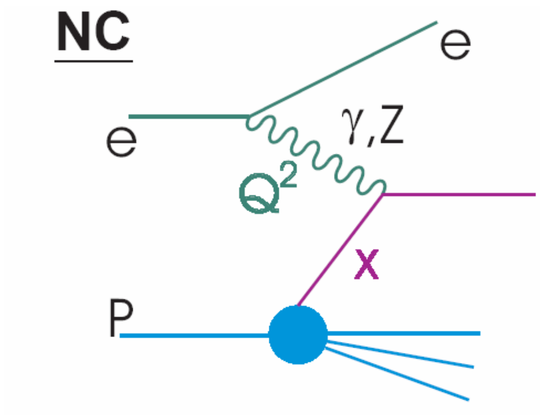
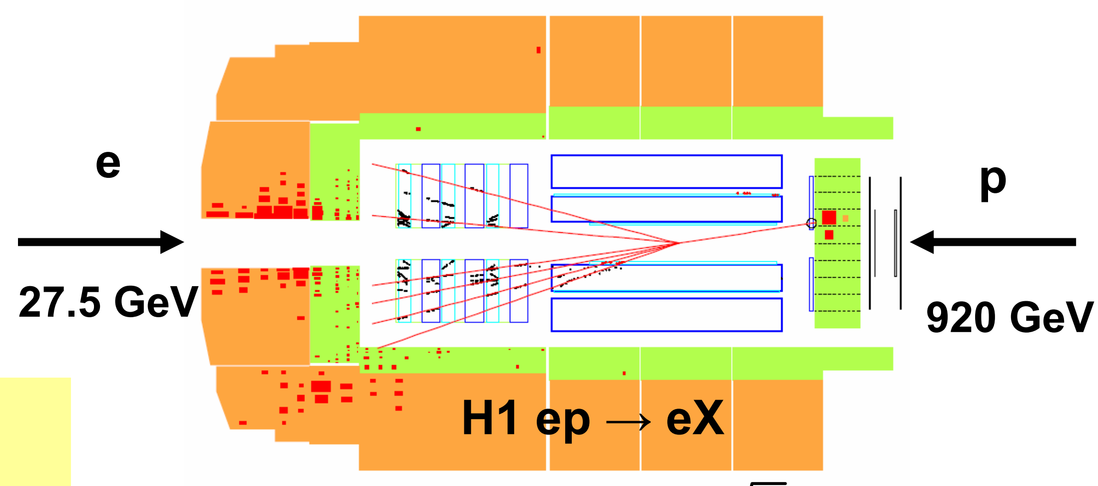
- QCD dynamics at low Bjorken–x
- QCD models
- Experimental Method
- Results
- Conclusions

Workshop on exclusive and diffractive processes at high energy proton-proton and nucleus-nucleus collisions, ECT* Trento, Feb 27 – March 2, 2012

Deep-inelastic ep scattering at HERA



DIS event at low Q^2



DIS phase space defined by :

Q^2 – virtuality of the exchanged boson

x – fraction of the proton momentum carried by the struck quark

$y = Q^2 / xs$ inelasticity

HERA : $\sqrt{s}$ = 319 GeV

precise measurement of the scattered electron (E_e scale at 1%) and hadronic final state (E_h scale at 1- 4 %)

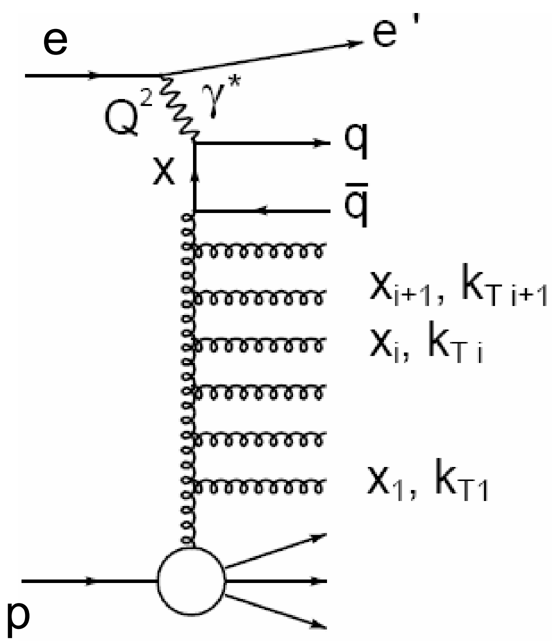
HERA : DIS at low Bjorken-x down to 10^{-5} → large γ*p centre-of-mass energy,

$W_{\gamma^*p} \approx Q^2 / x$

- **enhanced phase space for long parton cascades exchanged between the proton and the photon**

QCD dynamics at low Bjorken-x

pQCD – multiparton emissions described only with approximations



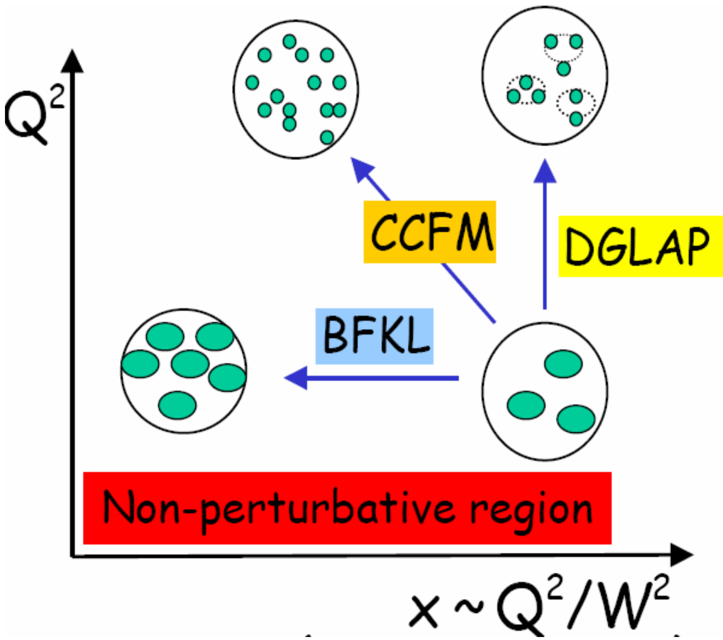
- **DGLAP : Dokshitzer-Gribov-Lipatov-Altarelli-Parisi evolution**
  applicable at large Q^2
  Emitted partons are ordered in k_T
  $k_{T,i}^2 \ll k_{T,i+1}^2 \ll \ldots \ll Q^2$
  Sums $(\alpha_S \ln Q^2)^n$ terms

- **BFKL : Balitsky-Fadin-Kuraev-Lipatov evolution scheme**
  at low x, in the Regge limit of pQCD $W^2_{\gamma^* p} \gg Q^2 \gg \Lambda^2_{QCD}$,
  transition from the DGLAP to BFKL approach expected
  No ordering in k_T, strong ordering in x_i
  Sums $(\alpha_S \ln(1/x))^n$ terms

- **CCFM : Ciafaloni-Catani-Fiorani-Marchesini equation**
  applicable at all x and Q^2
  Unification of the DGLAP and BFKL approaches
  Emitted partons are ordered in angles

QCD dynamics at low Bjorken-x

- Search at HERA for effects of parton dynamics beyond the standard DGLAP approach
- Define observables / phase space regions sensitive to low x effects

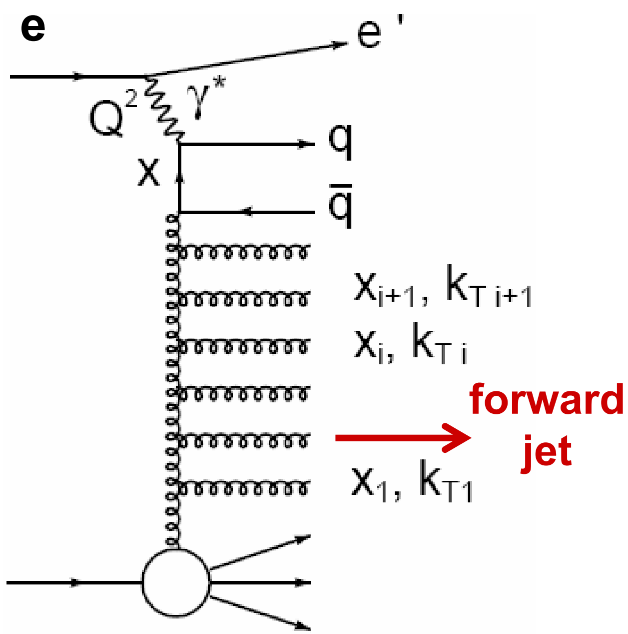


- Strong rise of the proton structure function $F_2(x, Q^2)$ with decreasing x
  – well described by NLO DGLAP over a large range of Q^2

F_2 measurement is too inclusive to discriminate between different QCD evolution schemes

Hadronic final states – reflect kinematics, structure of gluon emissions

(forward jets / particles, inclusive jets, multijet production, azimuthal correlation in dijet events, tranverse energy flow, pt distribution of hadrons)



BFKL - more hard partons emitted close to the proton

Mueller – Navelet jets in DIS (1990) :

High transverse momentum and high energy jets produced close to the proton remnant direction (forward region in LAB)

Suppress standard DGLAP evolution in Q^2 :

$$p^2_{T,jet} \approx Q^2$$

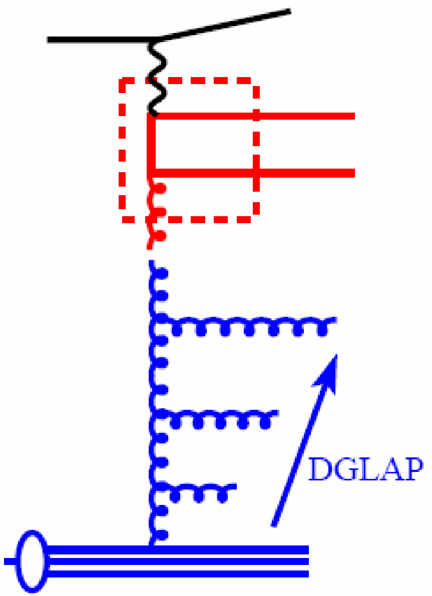
Enhance BFKL evolution in x :

$$x_{fwdjet} = E_{fwdjet} / E_p >> x_{Bj}$$

Monte Carlo models with different QCD dynamics

RAPGAP - DGLAP

LO QCD matrix elements + HO modelled by leading log parton showers

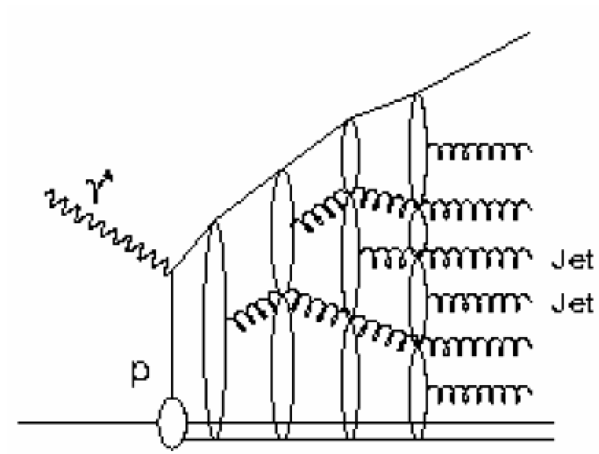


Single DGLAP ladder with strong ordering in k_T

ARIADNE Colour Dipole Model

CDM : QCD radiation from a colour dipole formed by the struck q and the p remnant.

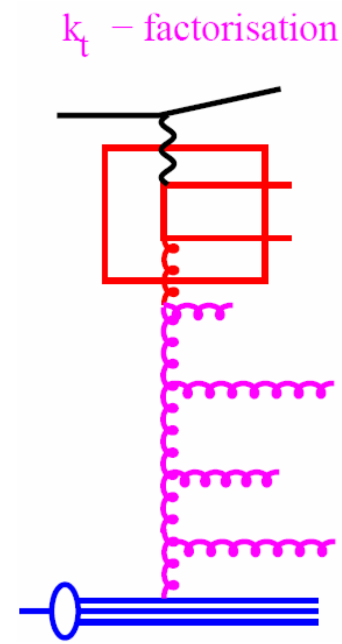
Chain of independently radiating dipoles formed by the emitted gluons.



BFKL-like Monte Carlo : random walk in k_T

CASCADE - CCFM

Off-shell QCD ME + parton emissions based on the CCFM equation.

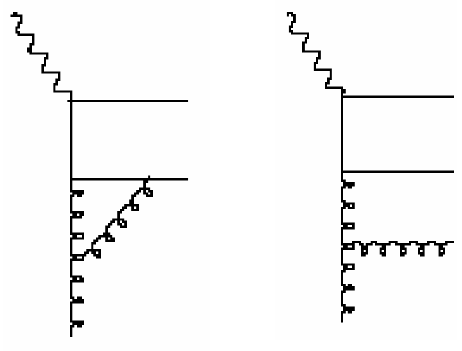


Input : unintegrated gluon density function, different uPDF sets include singular or full terms of the gluon splitting function.

Fixed order NLO DGLAP predictions

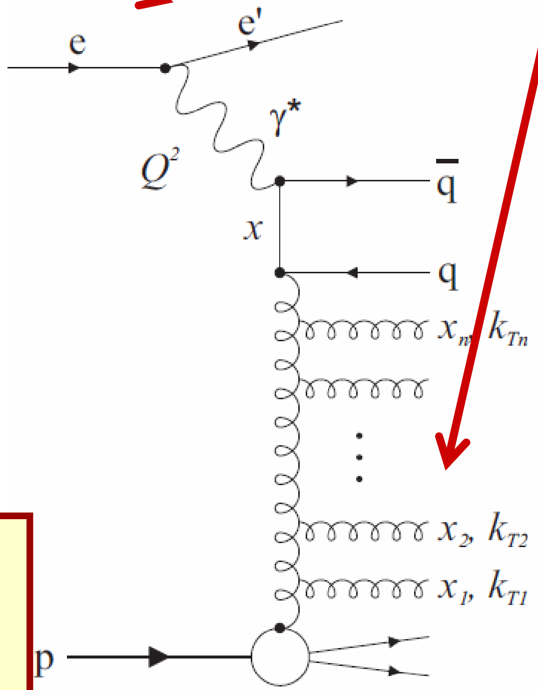
- Forward jet analysis – reconstruction of jets in the Breit frame
  → at least dijet topology

NLOJET ++ program (Nagy & Trocsanyi) :
dijet production at parton level in DIS at NLO (α_s^2) accuracy



- PDF : CTEQ6.6, $\alpha_s(M_Z) = 0.118$
- parton level cross sections corrected for hadronisation effects using the RAPGAP model

Azimuthal angle difference ΔΦ between the scattered positron and the forward jet may be sensitive to BFKL dynamics



- **Quark Parton Model e + q → e + q**

  simple two-body kinematics $\Delta\phi = \phi_{el} - \phi_{fwdjet} = \pi$

- **Inclusion of higher order processes decorrelates the jet from the positron**

As the rapidity distance approximated by $Y = \ln(x_{fwdjet}/x_{Bj})$ between the scattered positron and the forward jet grows the probability of multi-gluon emissions is increased

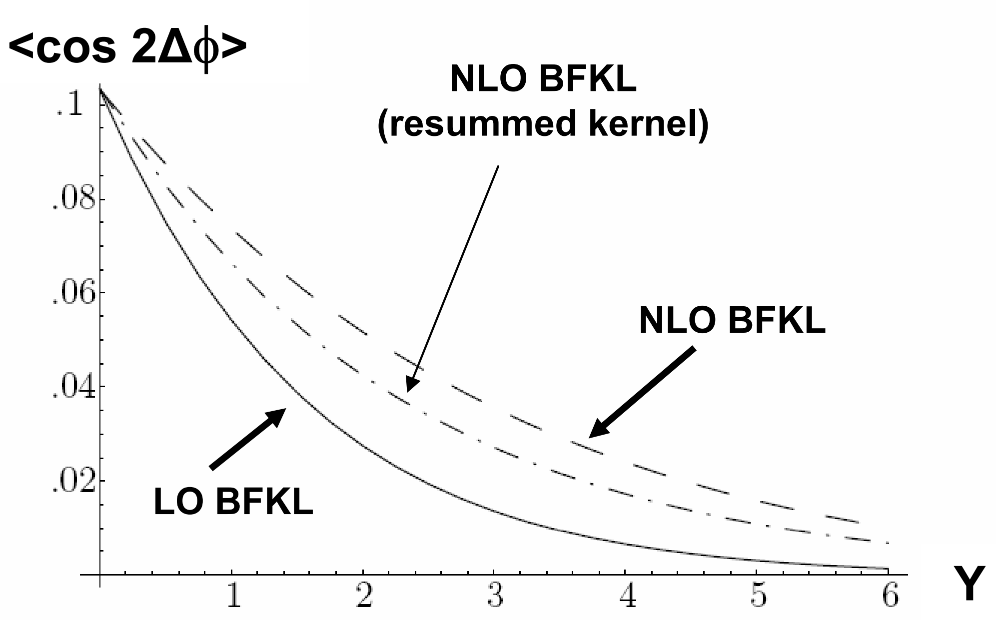
- **J. Bartels et al., Phys. Lett. B384(1996)300**
  calculation of Δϕ **distibution in LO BFKL**

- **S. Vera & F. Schwennsen, Phys. Rev. D77(2008)014001**
  calculation of Δϕ **distribution in NLO BFKL**

Forward jet production at NLO BFKL

S. Vera and F. Schwennsen, Phys. Rev. D77 (2008) 014001

BFKL kernel at NLO accuracy, jet vertex & photon impact factor using LO approximation



Results
for forward jets with ZEUS cuts

$20 < Q^2 < 100$ GeV^2
$0.05 < y < 0.7$
$4\cdot10^{-4} < x_{Bj} < 5\cdot10^{-3}$
$0.5 < p_t^2 / Q^2 < 2.0$

$Y = \ln(x_{jet} / x_{BJ})$ – evolution length in BFKL formalism

- The forward jet is more decorrelated from the scattered lepton for larger rapidity difference (center of mass energy)
- The azimuthal angle correlations increase when HO corrections are included for a fixed value of x_{Bj}
- Some angular deccorelation exits even for very small values of Y

Data selection

H1 experiment, HERA data (2000) with 38.2 pb^{-1}

DIS selection

$0.1 < y < 0.7$

$5 < Q^2 < 85\ GeV^2$

$0.0001 < x < 0.004$

Forward jets (inclusive k_T algorithm)

Jets reconstructed in the Breit frame and then boosted to LAB, all cuts in LAB

$p_{T,\ fwdjet} > 6\ GeV$

$1.73 < \eta_{fwdjet} < 2.79$

$x_{fwdjet} = E_{fwdjet} / E_p > 0.035$

$0.5 < p_{T,fwdjet}^2 / Q^2 < 6.0$

- suppress k_T ordered evolution by cut on p_T^2 / Q^2
- enhance phase space for BFKL evolution without k_T ordering by cut on x_{fwdjet}

~ 14000 DIS events with at least one forward jet

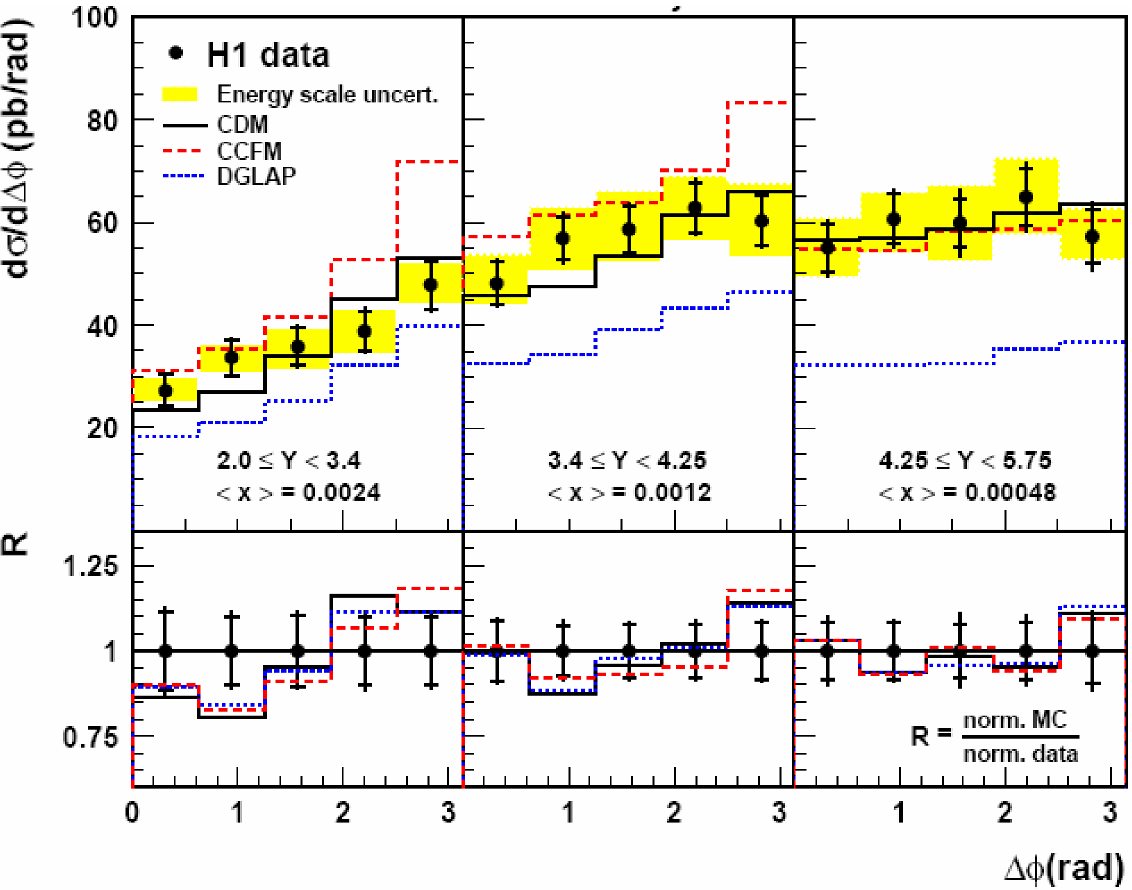
if more than one forward jet is found, the jet with the largest η_{fwdjet} is chosen

$\eta = -\ln(\tan \theta/2)$, θ defined with respect to the initial proton direction

Forward jet azimuthal correlations

At higher Y correspondig to lower x the forward jet is more decorrelated from the scattered electron

Positron – fwd jet rapidity distance $Y = \ln(x_{fwdjet} / x)$



Cross sections :

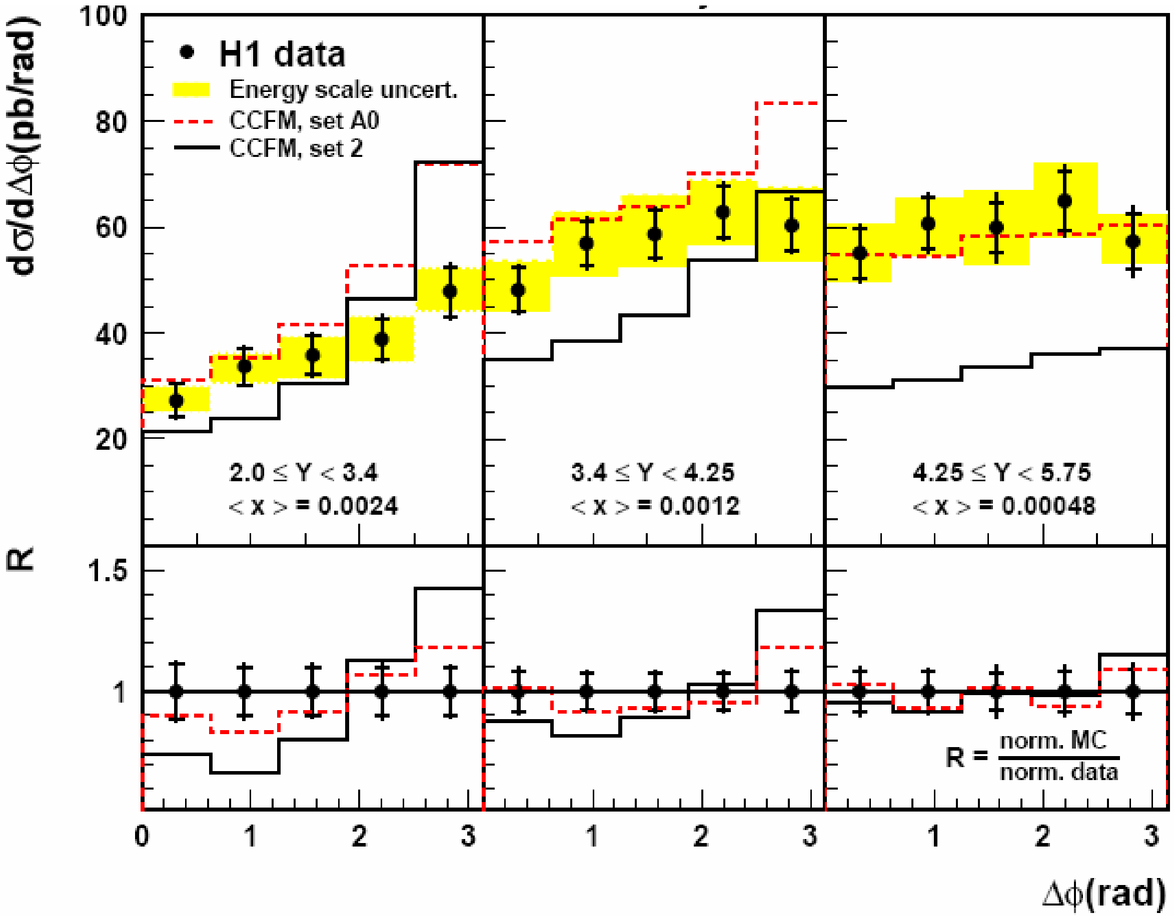
- well described by BFKL-like model CDM
- DGLAP predictions below the data
- CCFM as good description as CDM at large Y

Ratio R of MC to data for normalised cross-section

$$R = \left(\frac{1}{\sigma^{MC}} \frac{d\sigma^{MC}}{d\Delta\phi} \right) \Big/ \left(\frac{1}{\sigma^{data}} \frac{d\sigma^{data}}{d\Delta\phi} \right)$$

The shape of Δφ distributions is well described by all MC models

Predictions of the CCFM model depend on the choice of uPDF

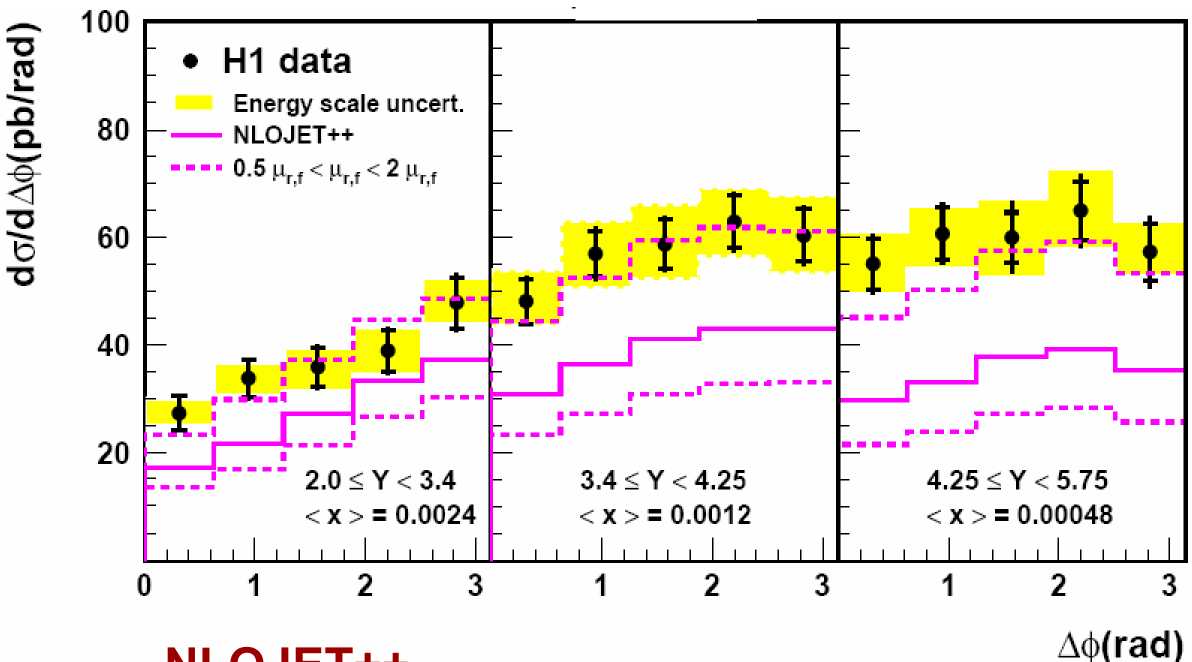


- **Cross sections** strongly depend on uPDF
- **Shape of Δφ distributions**
  - at low Y shows sensitivity to uPDF
  - well described by the set A0

Different splitting functions used in uPDF :

set A0 – only singular terms of the gluon splitting function
set 2 – includes also non-singular terms

Comparison to NLO ($O(\alpha_s^2)$)



NLO predictions

- shape of $\Delta\phi$ distributions described, but central value too low
- large scale uncertainty (of up to 50%) indicates importance of higher orders

NLOJET++

PDF : CTEQ6.6, $\alpha_S(M_Z)=0.118$

renormalisation and factorisation scales :

$$\mu_r^2 = \mu_f^2 = (p^2_{T,\,fwdjet} + Q^2)/2$$

theoretical uncertainty : factor 2 or ½ applied to μ_r and μ_f scales simultaneously

$\Delta\phi$ decorrelation :

no discrimination between different evolution schemes

- **forward jet originates from the hard matrix elements ?**

(similar in used MC models)

Studies of parton to hadron correlation with the DGLAP-based RAPGAP model :

► **define "distance measure" ΔR between parton jet and hadron jet**

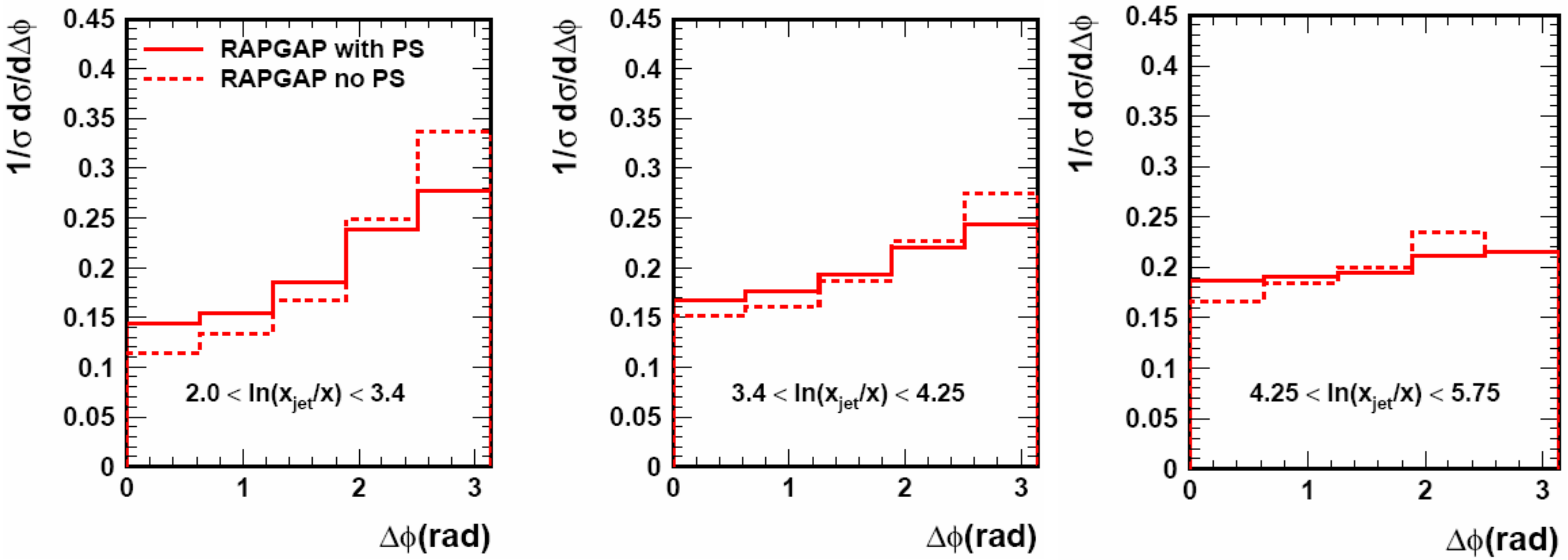
$$\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\Phi)^2}$$

$\Delta R < 0.5$ hadron jet is correlated to parton from ME / from parton shower

► **Y bin**	**forward jet originated from PS**
bin no. 1 ($2 < Y < 3.4$)	51.9%
bin no. 2 ($3.4 < Y < 4.25$)	67.5%
bin no. 3 ($4.25 < Y < 5.75$)	79.0%

Δϕ decorrelation :
why no dependence on parton shower ?

RAPGAP parton showers switched on / off



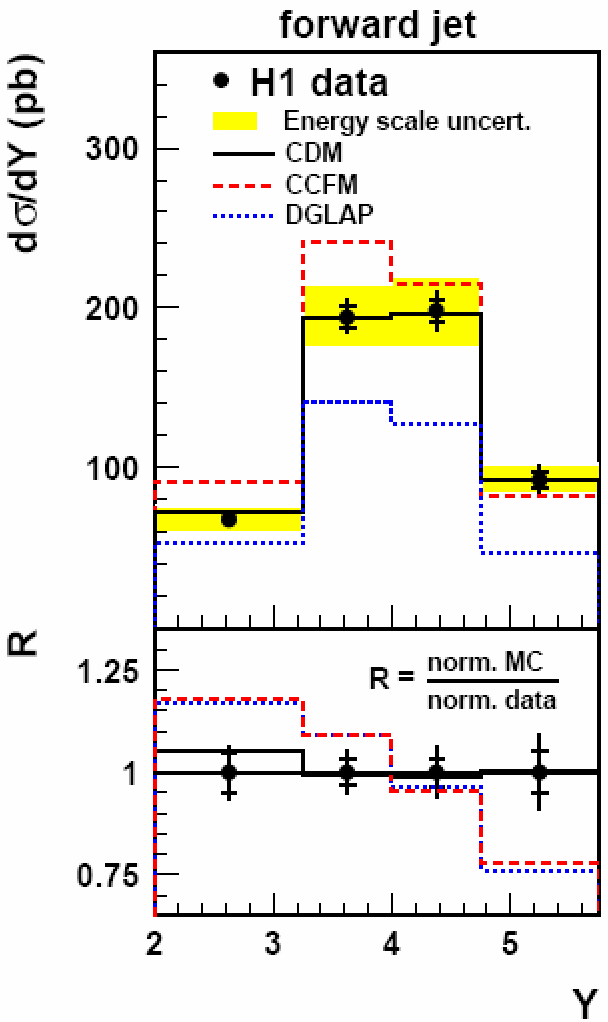
- The shape of the Δϕ distribution is only slightly changed when the initial state parton shower is switched off
- Δϕ decorrelation does not depend on Q^2 (not shown) - does **NOT** come from the photon virtuality

↓

Decorrelation in Δϕ is governed by the phase space requirements
(mainly by rapidity separation Y)

Normalisation of the cross sections is influenced by the amount of soft radiation from PS (depends on the evolution scheme)

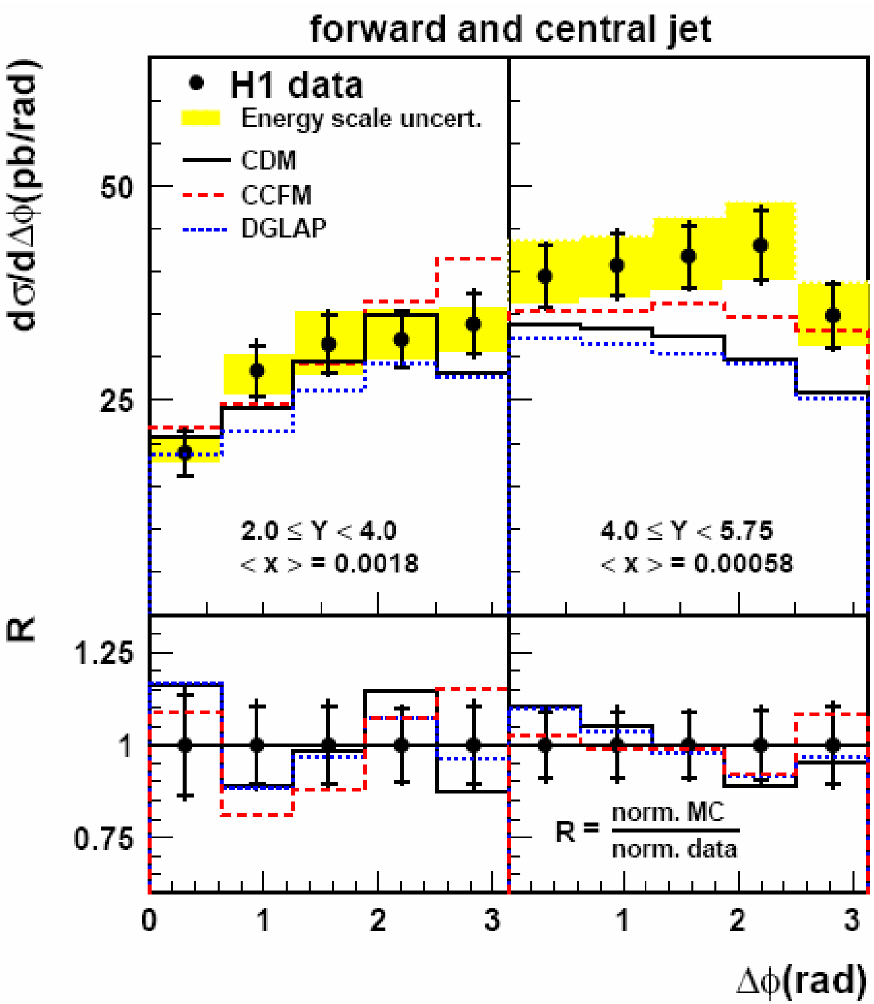
Forward jet cross section dσ / dY



- BFKL–like model CDM describes the data best
- DGLAP too low, especially at large Y
- CCFM predictions to high at low x, but describe the data at large Y

$Y = \ln(x_{jet} / x)$ rapidity separation between the most forward jet and the scattered positron

Forward and central jet cross sections $d\sigma / d\Delta\phi$

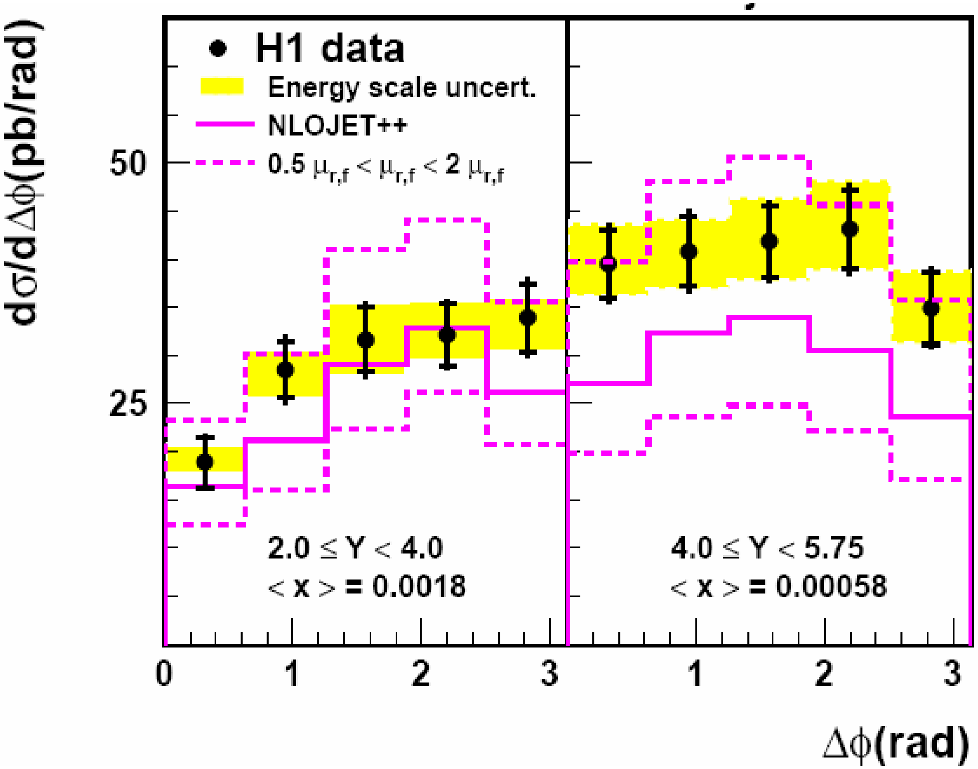


- Subsample of events with
  forward jet + additional central jet
  (~8900 events)
  $p_{T,cenjet} > 4$ GeV
  $-1 < \eta_{cenjet} < 1$
  $\Delta\eta = \eta_{fwdjet} - \eta_{cenjet} > 2$
  (enhance radiation between the forward and central jet)
- $\Delta\phi$ still between the forward jet and the scattered positron

- at low Y all models describe the data reasonably well
- at high Y all models are below the measurements
  - ► with CCFM closest to the data

Forward and central jet cross sections $d\sigma / d\Delta\phi$

Comparison to NLO ($O(\alpha_s^2)$)



NLO predictions

- at low Y reasonable description of the data
- at high Y, central value to small but still within theory uncertainty
- large scale uncertainty (of up to 40%) indicates importance of higher orders

forward jet + additional central jet sample

Conclusions

- Differential cross sections & normalised distributions have been measured as a function of $\Delta\phi$ and the rapidity separation Y, between the forward jet and the scattered positron
- Cross sections are best described by the BFKL-like model CDM
- DGLAP-based RAPGAP model is substantially below the data
- The CCFM model gives a reasonable description of the data but shows sizeable sensitivity to uPDF
- The shape of $\Delta\phi$ distributions is well described by MC models based on different QCD evolution schemes
- NLO DGLAP predictions are in general below the data, but still in agreement with the large theoretical uncertainties

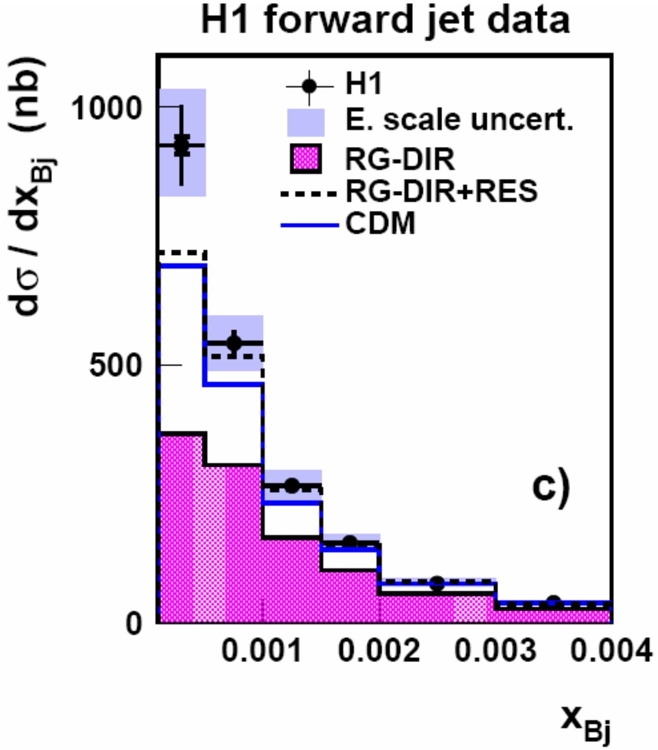
backup

Systematic uncertainties

	$d\sigma/d\Delta\phi$ fj	$d\sigma/d\Delta\phi$ fj + cj	$d\sigma/dY$
Model dependence (CDM,Rapgap)	2 – 6%		
LAr hadronic en. scale (±4%)	7 – 2%		
Spacal em en. scale (±1%)	below 3%		
Angle of scattered electron (± 1 mrad)	negligible effect		
Trigger	2 – 4%		
Luminosity	1.5%		
Total	11 – 12%		

Forward jet cross section dσ/dx

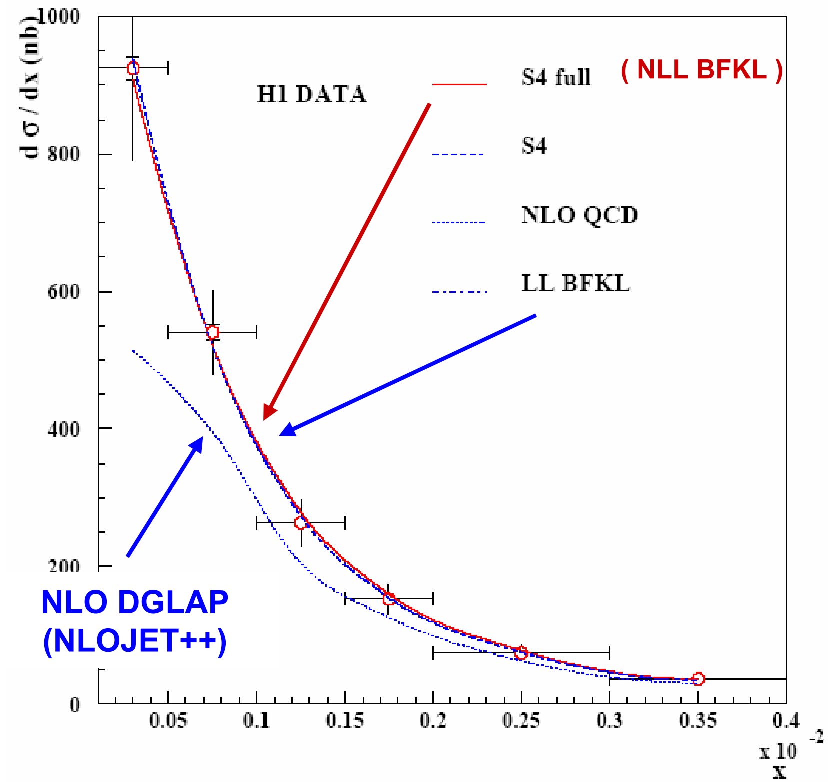
H1 data : Eur. Phys. J. C46 (2006)27



LO DGLAP (RG-DIR) below the data

CDM model and DGLAP resolved photon model (RG-DIR+RES) closest to the data, however the data are still below predictions at low x

BFKL calulations
Kepka, Royon, Marquet & Peschanski
Phys. Lett. B665 (2007) 236



NLO DGLAP below the data at low x

Difference between LL-BFKL and NLL-BFKL (NLL BFKL kernel + free normalisation parameter) is very small

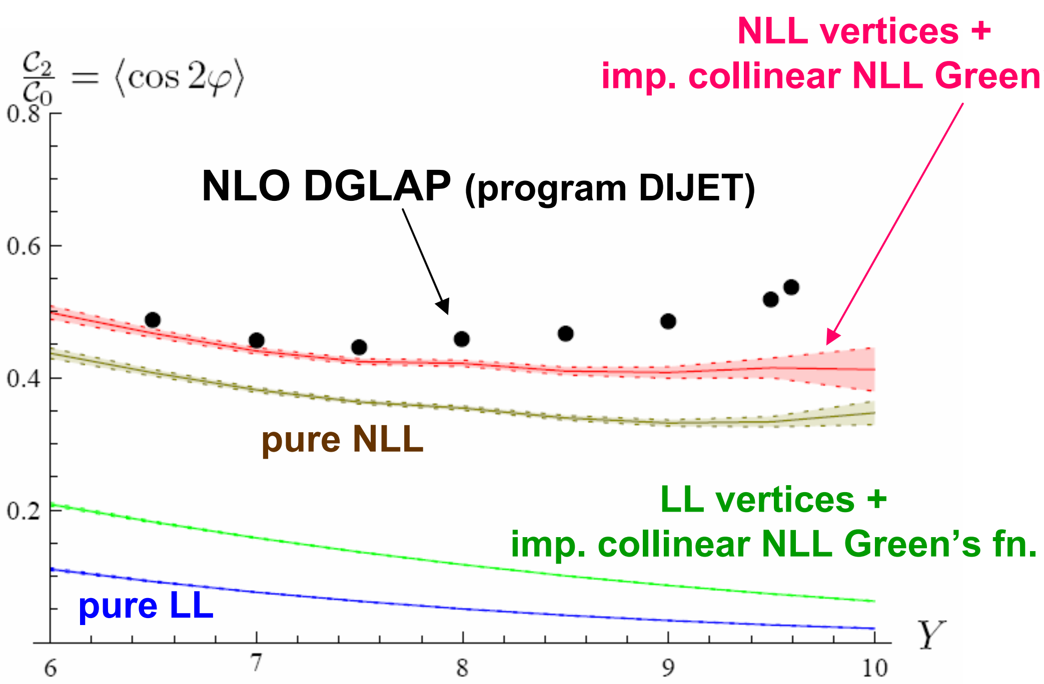
Mueller- Navelet jet at LHC – complete NLL BFKL calculations

Colferai, Schwennsen, Szymanowski & Wallon, JHEP 12(2010)026

next-to-leading corrections to the Green's function and to the Mueller-Navelet vertices

LHC $\sqrt{S}$ = 14 TeV, $p_{T,jet1}$ = 35 GeV, $p_{T,jet2}$ = 50 GeV

Azimuthal correlation $\langle \cos 2\phi \rangle = \langle \cos(2 \cdot (\phi_{jet1} - \phi_{jet2} - \pi)) \rangle$



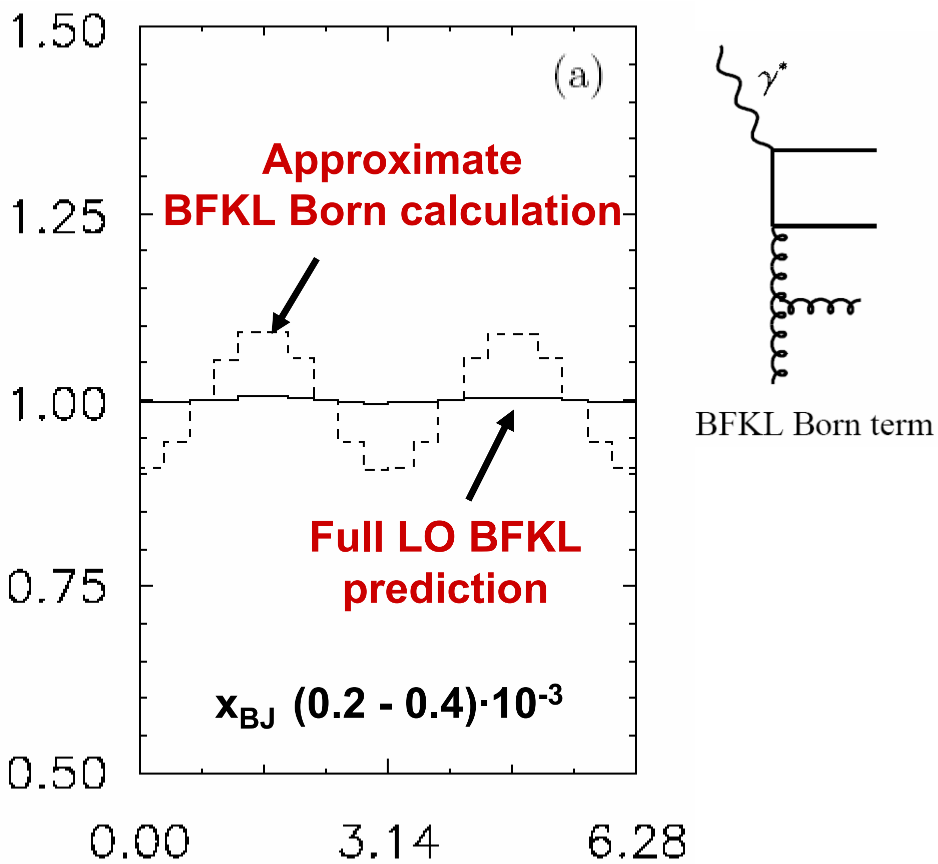
- **importance of NLL vertex corrections**
- **no significant difference between NLL BFKL and NLO DGLAP**

H1 measurements →
the electron-forward jet decorrelation in DIS does not discriminate between different evolution schemes

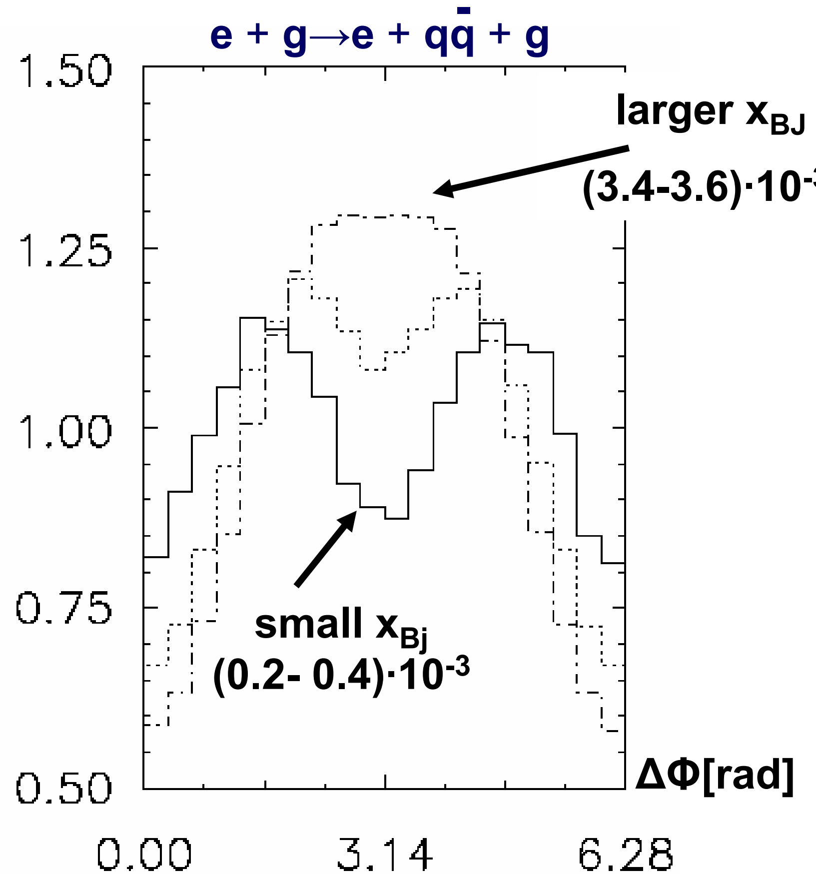
Forward jets in DIS, Bartels et al., 1996, $\Delta\Phi = \Phi_{el} - \Phi_{jet}$ in the LAB frame

normalised cross section

LO BFKL



Fixed order $O(\alpha_S^2)$ predictions



BFKL Born - clear maximum at $\Delta\Phi = \pi/2$
(BFKL ladder reduces to the gluon propagator)

Full LO BFKL – no Φ dependence

small x_{BJ} – fixed order $O(\alpha_S^2)$ and BFKL Born predictions are similar (max. at $\Delta\Phi \sim \pi/2$)